

Learning to Count up to Symmetry

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Abstract

In this paper we develop the theory of how to *count*, in thin concurrent games, the configurations of a strategy witnessing that it reaches a certain configuration of the game. This plays a central role in many recent developments in concurrent games, whenever one aims to relate concurrent strategies with weighted relational models.

The difficulty, of course, is symmetry: in the presence of symmetry many configurations of the strategy are, morally, different instances of the *same*, only differing on the inessential choice of copy indices. How do we know which ones to count? The purpose of the paper is to clarify that, uncovering many strange phenomena and fascinating pathological examples along the way.

To illustrate the results, we show that a collapse operation to a simple weighted relational model simply counting witnesses is preserved under composition, provided the strategies involved do not deadlock.

1 Introduction

Thin concurrent games [5] are a complex but powerful setting for truly concurrent game semantics; one of the latest iterations of a long line of work [1, 11, 13] on game semantics questioning the premise that a play should be a total chronological ordering. They are very expressive, able to express various languages both pure [4] and stateful [5]; including with various quantitative aspects [3, 7]. One strength of concurrent games in general is the clean link they offer with relational-like semantics: a strategy may (slightly naively) be seen as a collection of points of the *web* (in the sense of relational semantics) enriched with causal information. This enables a clean connection with the relational model, which served as basis *e.g.* for Melliès' fully complete model of linear logic [10] (see also [6]).

Now, relational semantics as well can be enriched with quantitative information; this is the basis for *probabilistic coherence spaces* [8]. Probabilistic coherence spaces are obtained via a biorthogonality construction on top of the relational model *weighted* by elements of $\overline{\mathbb{R}_+}$, the completion of non-negative reals \mathbb{R}_+ with a point at infinity. Instead of merely relations, morphisms from set A to set B are then matrices

$$(\alpha_{a,b})_{(a,b) \in A \times B} \in \overline{\mathbb{R}_+}^{A \times B}$$

composed via the potentially infinite matrix multiplication formula

$$(\beta \circ \alpha)_{a,c} = \sum_{b \in B} \alpha_{a,b} \cdot \beta_{b,c}. \quad (1)$$

Beyond real scalars, more generally one can construct a *weighted relational model* parametrized by certain semirings [9]. Adding typing information one goes beyond semirings, for instance the adequate model for the quantum λ -calculus of [12] uses weights from the category of finite dimensional Hilbert spaces and completely positive maps.

Above, we mentioned a *collapse* from concurrent games to the relational model. Does it hold with quantitative information? Such results appear in the literature [3, 7] – though we shall see in this paper that the definition of this collapse in [3] is not quite right. This seemingly simple question holds some surprises. This is the question that this paper solves; detailing the basis for part of [7], and identifying and correcting the mistake in [3].

As a matter of fact, the difficulty is not in handling the *weights*, but in *listing the right witnesses*: if (1) originates in a bijection between witnesses, then provided this bijection preserves the weights (and it will be generated in such a way that it does), it follows that adding weights is relatively painless. On the other hand, coming up with the right notion of witnesses is really hard. Indeed, in the presence of replication of resources, configurations in strategies are countably duplicated, so it is meaningless to sum over all of those as one does without symmetry. What are, then, the right witnesses? Symmetry classes of configurations? Something else? In this paper we give the answer, and illustrate it with a proof of a formula like (1) for a simple weighted relational model simply counting witnesses.

We shall see that the appealingly simple idea of [3] to use symmetry classes of configurations as witnesses is, in general, wrong. We give a more refined notion of witnesses, taking advantage of the split of the symmetry into *positive* and *negative* reindexings offered by thin concurrent games [5]. This lets us solve the problem, but with the cost of adding a new condition to thin concurrent games called *representability*, which states the existence, for every symmetry class, of a *canonical* representative on which the symmetry decomposes neatly into a positive and negative parts.

Outline. The structure of the paper is as follows. In Section 2 we fix the notations for thin concurrent games used in this paper and recall a few notions. In Section 3 we give a technical explanation of the problem and its difficulties. In Section 4 we introduce the new notions of *canonicity* and *representability*. In Section 5 we give the central contribution of the paper, the proof of (1). Finally, in Section 6 we give a few ending remarks.

2 Preliminaries

2.1 Notations and terminology

In this paper, we assume some familiarity with concurrent games, and more precisely with *thin concurrent games* [5]. Let us fix a few conventions for notations and terminology.

By *strategy* we will always mean \sim -strategy in the sense of [5]. We will sometimes refer to *pre- \sim -strategies*, which must be understood as in [5]. If $\sigma : S \rightarrow A^\perp \parallel B$ is a strategy from A to B , we write $\sigma : A \xrightarrow{S} B$. We often use x^S, y^S, \dots to range over configurations of S , with S as a superscript. If $x^S \in \mathcal{C}(S)$, we take the convention that

$$\sigma x^S = x_A^S \parallel x_B^S,$$

in the paper we will use $x_A^S \in \mathcal{C}(A)$ and $x_B^S \in \mathcal{C}(B)$ without further introduction.

If A is a tcg, we write \cong_A for its symmetry, and $\theta : x \cong_A y$ if the bijection $\theta : x \simeq y$ is in \cong_A – in which case we say that θ is a *symmetry*. For $x, y \in \mathcal{C}(A)$, we write $x \cong_A y$ for the induced equivalence relation. We use similar notations for the positive and negative sub-symmetries, with \cong_A^+ for the positive and \cong_A^- for the negative. We use for symmetries on strategies similar notations as for configurations. For $\sigma : A \xrightarrow{S} B$, we often tag symmetries in S with S , as in $\varphi^S : x^S \cong_S y^S$. Then, we write $\varphi_A^S : x_A^S \cong_A y_A^S$ and $\varphi_B^S : x_B^S \cong_B y_B^S$.

In diagrams, dotted lines signify immediate causal links in the game, whereas \rightarrow means immediate causality in the strategy. If the direction of causal links is unspecified (*e.g.* with dotted lines with no arrow head), then it must be read from top to bottom.

2.2 Interaction and composition

Consider two strategies $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$.

Recall that their *interaction*

$$\tau \circledast \sigma : T \circledast S \rightarrow A \parallel B \parallel C$$

has set $\mathcal{C}(T \circledast S)$ isomorphic to pairs $(x^S, x^T) \in \mathcal{C}(S) \times \mathcal{C}(T)$ such that $x_B^S = x_B^T = x_B$, and which are *causally compatible*, in the sense that the induced bijection

$$x^S \parallel x_C^T \simeq x_A^S \parallel x_B \parallel x_C^T \simeq x_A^S \parallel x^T$$

is secured [5]. We write $x^T \circledast x^S \in \mathcal{C}(T \circledast S)$ for the corresponding configuration; then:

$$(\tau \circledast \sigma)(x^T \circledast x^S) = x_A^S \parallel x_B \parallel x_C^T.$$

The composition

$$\tau \odot \sigma : A \xrightarrow{T \odot S} C$$

is obtained from the interaction through a hiding operation [5]. We recall:

Proposition 1. *The set $\mathcal{C}(T \odot S)$ is isomorphic to the set of pairs $(x^S, x^T) \in \mathcal{C}(S) \times \mathcal{C}(T)$ such that $x_B^S = x_B^T = x_B$, which are causally compatible and minimal, in the sense that if $y^S \subseteq x^S$ and $y^T \subseteq x^T$ are matching and causally compatible, and*

$$x_A^S \parallel x_C^T = y_A^S \parallel y_C^T,$$

then $x^S = y^S$ and $x^T = y^T$. If x^S and x^T are matching, causally compatible, and minimal, we write $x^T \odot x^S \in \mathcal{C}(T \odot S)$ for the corresponding configuration. We then have

$$(\tau \odot \sigma)(x^T \odot x^S) = x_A^S \parallel x_C^T.$$

Proof. Direct from the definition. If a pair (x^S, x^T) is matching and causally compatible, then it is minimal iff $x^T \otimes x^S$ has all its maximal events visible (*i.e.* in A or C); and those are in one-to-one correspondence with configurations of $T \odot S$. \square

Interaction behaves like a cartesian product (restricted to the matching causally compatible configurations), while composition has this additional minimality assumption. We wish to get rid of minimality, since we wish to link to weighted relational models, where (intuitively) a witness of the composition is a pair of witnesses. This can be achieved:

Definition 2. Let $\sigma : S \rightarrow A$ be a strategy.

A configuration $x \in \mathcal{C}(S)$ is $+$ -covered iff all its maximal events have positive polarity. We write $\mathcal{C}^+(S)$ for the set of $+$ -covered configurations of σ .

By extension, we say that $x^T \otimes x^S \in \mathcal{C}(T \otimes S)$ is $+$ -covered iff its maximal events are positive and write $x^T \otimes x^S \in \mathcal{C}^+(T \otimes S)$. This notion is useful, because we have:

Lemma 3. Consider $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ two strategies. Then, there is a bijection

$$\begin{aligned} \phi : \mathcal{C}^+(T \otimes S) &\simeq \mathcal{C}^+(T \odot S) \\ x^T \otimes x^S &\mapsto x^T \odot x^S \end{aligned}$$

such that if $(\tau \otimes \sigma)(x^T \otimes x^S) = x_A \parallel x_B \parallel x_C$, then $(\tau \odot \sigma)(\phi(x^T \otimes x^S)) = x_A \parallel x_C$.

Proof. If $x^T \otimes x^S \in \mathcal{C}^+(T \otimes S)$, then the pair (x^S, x^T) is automatically minimal: if not, then one can remove an event in B . But it must be negative for either σ or τ , contradiction. So we may simply set $\phi(x^T \otimes x^S) = x^T \odot x^S \in \mathcal{C}^+(T \odot S)$. \square

We have one last ingredient to introduce. One crucial difference between strategy composition and composition in weighted relational models, is that strategies may deadlock. This question is fairly well-explored; in particular in settings where we have performed such a collapse [3, 7, 6], we have done so under the assumption that strategies satisfied a condition called *visibility*, which prevents deadlocks [2]. Describing visibility is beyond the scope of this paper, but many of the results given here will be under the assumption that certain strategies do not deadlock. Accordingly, we define:

Definition 4. Strategies $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ do not deadlock iff for all $x^S \in \mathcal{C}(S)$, $x^T \in \mathcal{C}(T)$ and $\theta_B : x_B^S \cong_B x_B^T$, the composite bijection

$$x^S \parallel x_C^T \xrightarrow{\sigma \parallel x_C^T} x_A^S \parallel x_B^S \parallel x_C^T \xrightarrow{x_A^S \parallel \theta_B \parallel x_C^T} x_A^S \parallel x_B^T \parallel x_C^T \xrightarrow{x_A^S \parallel \tau^{-1}} x_A^S \parallel x^T$$

is secured.

This is, in particular, always the case when σ and τ are visible. If σ and τ do not deadlock then we may forget the causal compatibility condition in their interaction: configurations of the interaction correspond to arbitrary matching pairs.

We do *not* assume that all strategies considered do not deadlock. Throughout the paper, we make it explicit when we consider this hypothesis.

3 Towards a Quantitative Collapse

3.1 Relational collapse and symmetry

A game A has a natural associated notion of *position*, given by the set of *configurations* $\mathcal{C}(A)$. Configurations inform the relationship with relational-like semantics: if A is a game arising from a type in a *linear type system*, then the *web* (a set) interpreting this type in relational semantics may be identified with a subset of $\mathcal{C}(A)$ ¹. Likewise, a strategy

$$\sigma : A \xrightarrow{S} B$$

induces a relation $f\sigma = \{(x_A, x_B) \mid \exists x^S \in \mathcal{C}(S), \sigma x^S = x_A \parallel x_B\} \in \text{Rel}(\mathcal{C}(A), \mathcal{C}(B))$.

With this definition, for any $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ we automatically have that

$$f(\tau \odot \sigma) \subseteq (f\tau) \circ (f\sigma)$$

and the other inclusion holds if σ and τ do not deadlock.

This picture above is of course much simplified thanks to linearity. Without the linearity assumption, the games considered need to carry a *symmetry*. If A arises from a type, then the corresponding web is no longer (a subset of) $\mathcal{C}(A)$, but (a subset of) $\mathcal{C}_{\cong}(A)$, the set of *equivalence classes of configurations under symmetry*. In particular, we have

Lemma 5. *Consider N a negative tcg. Then,*

$$\mathcal{C}_{\cong}(!N) \cong \mathcal{M}_f(\mathcal{C}_{\cong}(N)).$$

where $\mathcal{M}_f(X)$ is the set of finite multisets of elements of set X .

Proof. Straightforward. □

We use $\mathbf{x}, \mathbf{y}, \dots$ as metavariables ranging over symmetry classes.

Above, $!$ stands for the AJM-style exponential described in Section 3.3.4 in [5]. Likewise, the reader familiar with relational semantics will recognize in $\mathcal{M}_f(X)$ the familiar exponential modality. This traces the path to extend the links between game and relational semantics beyond the linear case: simply correct the definition of $f\sigma$ by setting:

$$f\sigma = \{(\mathbf{x}_A, \mathbf{x}_B) \in \mathcal{C}_{\cong}(A) \times \mathcal{C}_{\cong}(B) \mid \exists x^S \in \mathcal{C}(S), x_A^S \in \mathbf{x}_A \ \& \ x_B^S \in \mathbf{x}_B\},$$

where $\sigma x^S = x_A^S \parallel x_B^S$, a naming convention that we shall adopt. If $x^S \in \mathcal{C}(S)$ is such that $x_A^S \in \mathbf{x}_A$ and $x_B^S \in \mathbf{x}_B$, we say that x^S is a **witness** for $(\mathbf{x}_A, \mathbf{x}_B)$ in $f\sigma$.

With this definition, it is immediate by definition of composition of strategies that we retain $f(\tau \odot \sigma) \subseteq f(\tau) \circ f(\sigma)$ for any strategies $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$.

¹Typically, in the presence of Question/Answer labeling, those are the *complete* configurations where every question is answered – but details do not matter for this paper.

3.2 Synchronization up to symmetry

More interesting is the reverse inclusion. Of course, the deadlock issue mentioned above still applies. But something else is also going on: consider $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$, and

$$(\mathbf{x}_A, \mathbf{x}_B) \in \int \sigma \quad (\mathbf{x}_B, \mathbf{x}_C) \in \int \tau.$$

By definition, this means that there are $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$ such that

$$x_A^S \in \mathbf{x}_A, \quad x_B^S \in \mathbf{x}_B, \quad x_B^T \in \mathbf{x}_B, \quad x_C^T \in \mathbf{x}_C.$$

In particular, since we have $x_B^S \in \mathbf{x}_B$ and $x_B^T \in \mathbf{x}_B$ it follows that there is a (non-unique)

$$\theta : x_B^S \cong_B x_B^T,$$

a symmetry on B . So the witnesses $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$ might not quite reach the same configuration of the game: typically, they might involve completely distinct copy indices, and θ carries a reindexing from one to the other. Independently of the deadlocks, if we wish to provide a witness $y \in \mathcal{C}(T \odot S)$ for $(\mathbf{x}_A, \mathbf{x}_C)$ in $\tau \odot \sigma$, we must in particular find some $y^S \in \mathcal{C}(S)$ and $y^T \in \mathcal{C}(T)$ such that

$$y_A^S \in \mathbf{x}_A, \quad y_B^S = y_B^T, \quad y_C^T \in \mathbf{x}_C,$$

matching on B on the nose. So starting from $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$, we must *reindex* them until they match on B on the nose. Of course, this issue already arises in the process of constructing a game semantics based on copy indices, to show that equivalence of (uniform) strategies up to the choice of copy indices is stable under composition.

In thin concurrent games, the main tool to deal with it is the *weak bipullback property*:

Lemma 6 (Weak bipullback property). *Let $\sigma : S \rightarrow A$ and $\tau : T \rightarrow A^\perp$ be pre- \sim -strategies. Let $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$ and $\theta : \sigma x^S \cong_A \tau x^T$, such that the composite bijection*

$$x^S \xrightarrow{\sigma} \sigma x^S \xrightarrow{\theta} \tau x^T \xrightarrow{\tau} x^T$$

is secured. Then, there are $y^S \in \mathcal{C}(S)$ and $y^T \in \mathcal{C}(T)$ causally compatible, $\theta^S : x^S \cong_S y^S$ and $\theta^T : y^T \cong_T x^T$, such that $\tau \theta^T \circ \sigma \theta^S = \theta$. Moreover, y^S, y^T are unique up to symmetry.

This appears as Lemma 3.23 in [5]. The intuition is that σ and τ play against each other, each replacing Player copy indices with one they are prepared to play. By \sim -receptivity, τ must be receptive to a change in copy indices made by Player, and reciprocally; so y^S and y^T may be constructed by induction on the causal structure induced by the securedness assumption. If $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$, and we have $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$ with

$$\theta : x_B^S \cong_B x_B^T,$$

we may apply the lemma above for $\sigma \parallel C^\perp \rightarrow A^\perp \parallel B \parallel C^\perp$ and $A \parallel \tau : A \parallel T \rightarrow A \parallel B^\perp \parallel C$. Provided some other argument ensures the securedness assumption, then we obtain

$$y^S \parallel y_C^T \in \mathcal{C}(S \parallel C) \quad y_A^S \parallel y^T \in \mathcal{C}(A \parallel T)$$

matching on B ; and so we have found an interaction

$$y^T \circledast y^S \in \mathcal{C}(T \circledast S)$$

with $(\tau \circledast \sigma)(y^T \circledast y^S) = y_A^S \parallel y_B \parallel y_C^T$, satisfying $y_A^S \in \mathbf{x}_A$ and $y_C^T \in \mathbf{x}_C$ thus providing through hiding the desired witness for $(\mathbf{x}_A, \mathbf{x}_C) \in \int(\tau \odot \sigma)$.

3.3 Quantitative extension

But the above is purely qualitative: if $\sigma : A \xrightarrow{S} B$ then the collapse above lets us define which pairs $(\mathbf{x}_A, \mathbf{x}_B)$ are “inhabited” by σ . This is sufficient in order to link game semantics with relational semantics. But this is not sufficient if we want to reproduce this feat in the presence of *quantitative* information, such as probabilities or quantum valuations.

For the purposes of this paper, let us say that we are now interested not in the *mere existence* of a witness $x^S \in \mathcal{C}(S)$ such that $x_A^S \in \mathbf{x}_A$ and $x_B^S \in \mathbf{x}_B$, but in *counting* such witnesses. For reasons explained in Section 2.2, from now on we consider witnesses for $(\mathbf{x}_A, \mathbf{x}_B)$ not merely those configurations $x^S \in \mathcal{C}(S)$ such that $x_A^S \in \mathbf{x}_A$ and $x_B^S \in \mathbf{x}_B$; but those that are additionally $+$ -covered, *i.e.* we have $x^S \in \mathcal{C}^+(S)$.

From a strategy $\sigma : A \xrightarrow{S} B$, we want a $\overline{\mathbb{N}}$ -weighted relation, *i.e.* a function

$$\int \sigma : \mathcal{C}_\cong(A) \times \mathcal{C}_\cong(B) \rightarrow \overline{\mathbb{N}},$$

where $\overline{\mathbb{N}} = \mathbb{N} \cup \{+\infty\}$, *counting* the number of distinct witnesses for $(\mathbf{x}_A, \mathbf{x}_B)$. In that case, for $\mathbf{x}_A \in \mathcal{C}_\cong(A)$ and $\mathbf{x}_B \in \mathcal{C}_\cong(B)$, write $(\int \sigma)_{\mathbf{x}_A, \mathbf{x}_B} \in \overline{\mathbb{N}}$ for the corresponding coefficient.

In the spirit of weighted relations [9], we then want to prove that for all $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ that do not deadlock, we have that for all $\mathbf{x}_A \in \mathcal{C}_\cong(A)$ and $\mathbf{x}_C \in \mathcal{C}_\cong(C)$,

$$(\int(\tau \odot \sigma))_{\mathbf{x}_A, \mathbf{x}_C} = \sum_{\mathbf{x}_B \in \mathcal{C}_\cong(B)} (\int \sigma)_{\mathbf{x}_A, \mathbf{x}_B} \times (\int \tau)_{\mathbf{x}_B, \mathbf{x}_C}. \quad (2)$$

The convergence of the sum on the right hand side is ensured by the fact that we consider the completed natural numbers $\mathbb{N} \cup \{+\infty\}$ as in the weighted relational model.

How might we, from $\sigma : A \xrightarrow{S} B$, extract the weighted relation $\int \sigma$? Intuitively, we need

$$(\int \sigma)_{\mathbf{x}_A, \mathbf{x}_B} = |\mathbf{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)|$$

where $\mathbf{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)$ captures the *witnesses* in σ for symmetry classes $\mathbf{x}_A \in \mathcal{C}_\cong(A)$ and $\mathbf{x}_B \in \mathcal{C}_\cong(B)$, and where $|X|$ simply computes the cardinal, taken to be $+\infty$ for X infinite.

Situations where strategies carry additional weights, say probabilities or quantum valuations, would be dealt with similarly. In any case, the first obstacle to overcome is then to give a satisfactory definition of $\text{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)$.

Of course counting all $x^S \in \mathcal{C}^+(S)$ such that $x_A^S \in \mathbf{x}_A$ and $x_B^S \in \mathbf{x}_B$ makes no sense: there are almost always infinitely many of them since *e.g.* the construction $!N$ introduces countably many copy indices. The definition of witnesses must take symmetry into account.

Symmetry classes. The obvious candidate for witnesses, chosen in [3], is:

$$\text{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B) = \{x^S \in \mathcal{C}_\cong^+(S) \mid \forall x^S \in \mathbf{x}^S, x_A^S \in \mathbf{x}_A \ \& \ x_B^S \in \mathbf{x}_B\},$$

i.e. the symmetry classes of $+$ -covered configurations mapping to $\mathbf{x}_A \parallel \mathbf{x}_B$. This convincingly simple definition in fact hides a major subtlety. Indeed, (2) hints at a bijection

$$\text{wit}_{\tau \odot \sigma}(\mathbf{x}_A, \mathbf{x}_C) \cong \sum_{\mathbf{x}_B \in \mathcal{C}_\cong(B)} \text{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B) \times \text{wit}_\tau(\mathbf{x}_B, \mathbf{x}_C).$$

This seems straightforward. Firstly, if $\mathbf{z} \in \text{wit}_{\tau \odot \sigma}(\mathbf{x}_A, \mathbf{x}_C)$, then any choice $z \in \mathbf{z}$ is $z = z^T \odot z^S \in \mathcal{C}(T \odot S)$ and the symmetry classes of its projections yield

$$z^S \in \text{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B), \quad z^T \in \text{wit}_\tau(\mathbf{x}_B, \mathbf{x}_C),$$

for some $\mathbf{x}_B \in \mathcal{C}_\cong(B)$. These data are easily shown to be invariant under the choice of z .

Reciprocally, if $x^S \in \text{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)$ and $x^T \in \text{wit}_\tau(\mathbf{x}_B, \mathbf{x}_C)$, we may take arbitrary $x^S \in \mathbf{x}^S, x^T \in \mathbf{x}^T$, and via Lemma 6 find symmetric $y^S \in \mathbf{x}^S$ and $y^T \in \mathbf{x}^T$ agreeing on B on the nose. We may then form $y^T \odot y^S \in \mathcal{C}^+(T \odot S)$ and take its symmetry class in $\text{wit}_{\tau \odot \sigma}(\mathbf{x}_A, \mathbf{x}_C)$.

But one should not skip the details²: we must show that this construction only depends on the symmetry classes \mathbf{x}^S and \mathbf{x}^T , not on the specific choices $x^S \in \mathbf{x}^S$ and $x^T \in \mathbf{x}^T$ and the symmetry $\theta_B : x_B^S \cong_B x_B^T$ used to link them. But surely, that must be true, right?

Well, about that... It certainly was a surprise to us that the symmetry class obtained through synchronization *does* depend on the symmetry θ_B .

Example 7. Consider the following games. Firstly, $A = \emptyset$ is the empty game. Secondly, $C = (!\ominus)^\perp$ which has countably many Player moves written \checkmark_i for all $i \in \mathbb{N}$, all symmetric – we adopt here a convention followed throughout the paper: copy indices appear in grey, to distinguish them from other indices.

Thirdly, consider the game $B = !_{HO}(\ominus \rightarrow \oplus)$, where $!_{HO}$ is the “HO exponential” defined in Definition 2.24 with symmetries in Definition 2.27 in [5] (see also Proposition 3.3). This game has events, polarities and causal dependency those pictured in:



²We were guilty of that in [3].

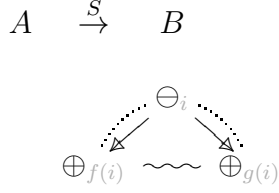


Figure 1: $\sigma : A \xrightarrow{S} B$

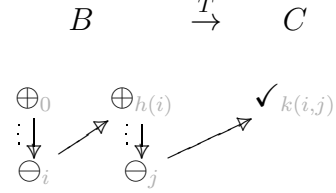
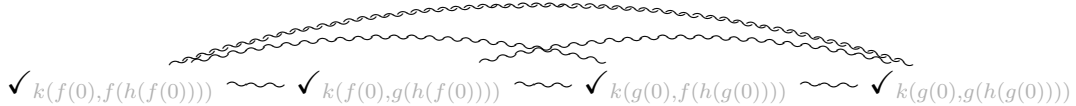


Figure 2: $\tau : B \xrightarrow{T} C$

with all finite sets consistent. Its symmetry comprises all order-isomorphisms between configurations. Its positive symmetry comprises all order-isomorphisms that preserve the initial (negative) move. Its negative symmetry comprises all order-isomorphisms such that $\theta(\oplus_{i,j}) = \oplus_{i',j}$ for some $i' \in \mathbb{N}$, *i.e.* they preserve the j component of the positive move. In practice, we will omit the first copy index for the event in the second row, which is redundant with the immediate causal antecedent of the event.

We now introduce two strategies $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ that we wish to compose, represented on Figures 1 and 2 where the functions f, g, h , and k are assumed injective, and 0 is not in the codomain of h . Note that the representation is symbolic: the diagrams must be understood by stating that every positive move has one copy for each instantiation of the metavariables $i, j \in \mathbb{N}$, with dependencies as indicated in the diagram. These copies are compatible with each other. Finally, the symmetries comprise order-isomorphisms that differ only by the value of the metavariables $i, j \in \mathbb{N}$. In particular, the two moves in the conflicting branches of σ are *not* symmetric (that would anyway contradict *thinness*).

First, we compute the composition $\tau \odot \sigma$, and observe that it is:



There are four events, pairwise conflicting, reflecting the two non-deterministic choices arising from the two calls to σ – one can read back which non-deterministic choice gave rise to which result from the copy indices, but that is another story. None of these events are symmetric: again, this would contradict *thinness*.

Now, let us define two configurations $x^S \in \mathcal{C}(S)$ and $x^T \in \mathcal{C}(T)$ as

$$x^S = \begin{array}{cc} \ominus_0 & \ominus_{h(f(0))} \\ \vdots & \vdots \\ \oplus_{f(0)} & \oplus_{g(h(f(0)))} \end{array} \quad x^T = \begin{array}{cc} \oplus_0 & \oplus_{h(f(0))} \\ \vdots & \vdots \\ \ominus_{f(0)} & \ominus_{g(h(f(0)))} \end{array} \parallel \checkmark_{k(f(0), g(h(f(0))))}$$

These two configurations match on B (and are causally compatible); and their composition yields the configuration $\{\checkmark_{k(f(0), g(h(f(0))))}\}$. We of course obtain the same result if we synchronize them through the trivial symmetry on their common interface:

$$\text{id} : x_B^T \cong_B x_B^T.$$

But there is another endosymmetry on $x_B = x_B^S = x_B^T$, namely

$$\mathbf{sw} : x_B^T \cong_B x_B^S,$$

exchanging the two copies. Synchronizing x^S and x^T through \mathbf{sw} via Lemma 6 instead gives:

$$\{\checkmark_{k(g(0), f(h(g(0))))}\}$$

which is not symmetric to $\{\checkmark_{k(f(0), g(h(f(0))))}\}$ in $T \odot S$. Indeed, intuitively, in x^S we only have the information that there were two calls to σ , with distinct non-deterministic resolutions. We do not know, just by looking at x^S , which one is the “first call” and which one is the “second call”. The symmetry $\theta : x_B^S \cong_B x_B^T$ “plugs” the two calls in x^T to their two non-deterministic resolutions in x^S . With id the first call selected $\oplus_{f(i)}$ and the second call $\oplus_{g(i)}$, and the other way around for \mathbf{sw} ; leading to non-symmetric outcomes.

Well, this is puzzling. If the obvious candidate for a bijection between witnesses $z \in \mathbf{wit}_{\tau \odot \sigma}(\mathbf{x}_A, \mathbf{x}_C)$ and pairs of witnesses $z^S \in \mathbf{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)$ and $z^T \in \mathbf{wit}_\tau(\mathbf{x}_B, \mathbf{x}_C)$ for some \mathbf{x}_B does not work, how can we hope to obtain (2)? This makes one wonder by what miracle the weighted relational model works at all – what does it really count?

Concrete witnesses. To investigate this issue we introduce an alternative, more concrete choice for witnesses. It is rooted in the following fact (Lemma 3.28 in [5]):

Lemma 8. *Let $\sigma : S \rightarrow A$ be a pre- \sim -strategy on A , and let $\theta : x \cong_S y$ such that $\sigma\theta \in \cong_A^+$. Then, $x = y$ and $\theta = \text{id}_x$.*

For this, the condition *thin* plays a crucial role. Intuitively, *thinness* means that the strategy has a canonical choice of copy indices for its moves, once Opponent fixes their choice of copy indices. Accordingly, the lemma above may be interpreted as saying that provided we remain in the positive symmetry (*i.e.* we do not change Opponent’s copy indices), then the choice of the *concrete* configuration $x \in \mathcal{C}(S)$ is unique. This suggests that we might take $\mathbf{wit}_\sigma(\mathbf{x}_A, \mathbf{x}_B)$ to range over concrete configurations of S matching with the game up to *positive* symmetry – of course, for that we need reference concrete configurations of the game rather than symmetry classes. So let us fix a choice, for any tcg A and any symmetry class $\mathbf{x}_A \in \mathcal{C}_\cong(A)$, of a concrete representative written $\underline{\mathbf{x}}_A \in \mathbf{x}_A$.

Our alternative definition of witnesses is, for $\sigma : A \xrightarrow{S} B$, $\mathbf{x}_A \in \mathcal{C}_\cong(A)$ and $\mathbf{x}_B \in \mathcal{C}_\cong(B)$:

$$\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) = \{x^S \in \mathcal{C}^+(S) \mid x_A^S \cong_A^- \underline{\mathbf{x}}_A \ \& \ x_B^S \cong_B^+ \underline{\mathbf{x}}_B\}$$

It will turn out (see Section 6) that (even assuming representability) these two notions of witnesses are *not* equivalent: the weighted relational model counts not symmetry classes, but concrete witnesses up to positive symmetry. In the rest of this paper, we aim to prove (as mentioned above, modulo one additional condition on games) that \mathbf{wit}^+ , unlike \mathbf{wit} , does the trick³. This will be quite the ride, so switch off your phone, fasten your seat belt, as we must now embark on a journey into the darkest corners of thin concurrent games.

³An early sign that \mathbf{wit}^+ is better behaved is that unlike \mathbf{wit} , it does not depend on the choice of the symmetry for σ – recall from Section A.1.2 in [5] that the symmetry is *not* unique.

4 Canonical configurations and representable games

4.1 Canonical representatives of symmetry classes

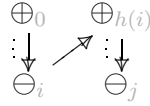
To motivate the development of this section, let us look at the definition just above:

$$\text{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) = \{x^S \in \mathcal{C}^+(S) \mid x_A^S \cong_A^- \mathbf{x}_A \ \& \ x_B^S \cong_B^+ \mathbf{x}_B\}.$$

This definition depends on a choice of a representative \mathbf{x}_A , once and for all, for every symmetry class \mathbf{x}_A . Of course, the set of witnesses we obtain this way depends on this choice: a different choice of representatives yields configurations of S where Opponent uses different copy indices. But what we really need for this definition to be of any use, is that the *cardinal* of $\text{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)$ should not depend on the representatives $\mathbf{x}_A, \mathbf{x}_B$.

Bad news: it does.

Example 9. Remember the game $B = !_{HO}(\ominus \rightarrow \oplus)$ of Example 7. Consider the strategy



written $\sigma : S \rightarrow B^\perp$, which is τ of Example 7 without the last move.

Now, imagine that we fix as representative for a symmetry class in B^\perp the configuration:

$$\mathbf{x}_B = \begin{array}{cc} \oplus_1 & \oplus_2 \\ \vdots & \vdots \\ \ominus_1 & \ominus_2 \end{array}.$$

Let us consider the configurations of S matching \mathbf{x}_B up to positive symmetry. First, a configuration $x \in \mathcal{C}(S)$ matching our requirements has four moves, and each Player move has *exactly one* successor. So it must have the following form, for some $i, j \in \mathbb{N}$,

$$\begin{array}{cc} \oplus_1 & \oplus_{h(i)} \\ \vdots & \vdots \\ \ominus_i & \ominus_j \end{array} \quad (3)$$

and finding the witnesses for \mathbf{x}_B boils down to figuring out all possible positive symmetries

$$\theta : \begin{array}{cc} \oplus_1 & \oplus_{h(i)} \\ \vdots & \vdots \\ \ominus_i & \ominus_j \end{array} \cong_{B^\perp}^+ \begin{array}{cc} \oplus_1 & \oplus_2 \\ \vdots & \vdots \\ \ominus_1 & \ominus_2 \end{array}$$

The positive symmetry of B^\perp is the negative symmetry of B : it lets us change the indices of minimal events, but the second component for positive events must be left unchanged. We may freely associate the minimal events either as $\oplus_1 \leftrightarrow \oplus_1$ and $\oplus_{h(i)} \leftrightarrow \oplus_2$; or as $\oplus_1 \leftrightarrow \oplus_2$ and $\oplus_{h(i)} \leftrightarrow \oplus_1$. But if we do the former, as the symmetry is positive it

forces $i = 1$ and $j = 2$. Likewise, if we do the latter, it forces $i = 2$ and $j = 1$. So overall, there are *exactly two* configurations of S matching \underline{x}_B up to positive symmetry:

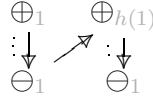


In particular, there are *two* witnesses for \underline{x}_B . This is confusing, because these two configurations are symmetric in S , so we seem to be counting the same symmetry class of S twice – and we shall see indeed that this is a pathological example.

In contrast, assume we pick as representative for \underline{x}_B the following configuration:

$$\underline{x}'_B = \begin{array}{c} \oplus_1 \quad \oplus_2 \\ \vdots \quad \vdots \\ \ominus_1 \quad \ominus_1 \end{array}.$$

Now, there is only *exactly one* configuration of S matching \underline{x}'_B up to positive symmetry:



Indeed, starting from (3), the positive symmetry forces i and j to be both 1; and we obtain the unique configuration above. So the choice of \underline{x}_B affects the number of witnesses.

What is the moral of the story? This is subtle. Notice that while there is indeed *exactly one* configuration $x \in \mathcal{C}(S)$ matching \underline{x}'_B up to positive symmetry, there are still *two symmetries* $\theta : \sigma x \cong_{B^\perp}^+ \underline{x}'_B$, corresponding to $\{\oplus_1 \leftrightarrow \oplus_1, \oplus_{h(1)} \leftrightarrow \oplus_2\}$ and $\{\oplus_1 \leftrightarrow \oplus_2, \oplus_{h(1)} \leftrightarrow \oplus_1\}$. So for \underline{x}_B we get *two* witnesses, and each has *one* positive symmetry to \underline{x}_B ; while for \underline{x}'_B we get *one* witness, with *two* positive symmetries. So the mismatch between the representatives is explained if one factors in the number of positive symmetries.

To comment further: there are two positive endo-symmetries $\underline{x}'_B \cong_{B^\perp}^+ \underline{x}'_B$: the identity, and the swap between positive events. In contrast, in \underline{x}_B , swapping the positive events

$$\begin{array}{c} \oplus_1 \quad \oplus_2 \\ \vdots \quad \vdots \\ \ominus_1 \quad \ominus_2 \end{array} \cong_{B^\perp}^+ \begin{array}{c} \oplus_2 \quad \oplus_1 \\ \vdots \quad \vdots \\ \ominus_1 \quad \ominus_2 \end{array}$$

while preserving Opponent indices cannot be achieved via an endosymmetry, this requires changing the configuration. To avoid such pathological cases, we must select \underline{x}_B such that the positive symmetry whose effect is, intuitively, merely to swap (the copy indices of) two Player events, still has \underline{x}_B as codomain. We do not have a definition capturing exactly this, as it is not clear how to formalize this idea of the minimal symmetry “swapping two Player events”. However, for our purposes the following definition does the job.

Definition 10. Consider A a tcg, and $x \in \mathcal{C}(A)$.

We say that x is **canonical** iff any $\theta : x \cong_A x$ factors uniquely as

$$x \stackrel{\theta^-}{\cong}_A^- x \stackrel{\theta_+}{\cong}_A^+ x,$$

with in particular x in the middle.

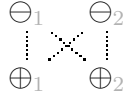
So endo-symmetries of canonical configurations decompose as endo-symmetries, positive and negative. Of course we already know that all endosymmetries (like all symmetries) decompose as the composite of a positive and a negative symmetries (see Lemma 3.19 of [5]). But there is a priori no reason why the decomposition should have the same configuration in the middle. This is in fact not always the case: for instance, picking the problematic configuration \underline{x}_B of the example above, we have the decomposition

$$\begin{array}{ccc} \begin{array}{c} \oplus_1 \\ \vdots \\ \ominus_1 \end{array} & \begin{array}{c} \oplus_2 \\ \vdots \\ \ominus_2 \end{array} & \cong_{B^\perp}^- \\ \begin{array}{c} \oplus_1 \\ \vdots \\ \ominus_2 \end{array} & \begin{array}{c} \oplus_2 \\ \vdots \\ \ominus_1 \end{array} & \cong_{B^\perp}^+ \\ \begin{array}{c} \oplus_2 \\ \vdots \\ \ominus_2 \end{array} & \begin{array}{c} \oplus_1 \\ \vdots \\ \ominus_1 \end{array} & \end{array}$$

where rather than drawing the symmetries, we suggest them by considering that they preserve the position of events in the diagrams. If we wish to avoid the problem mentioned above, we must project strategies only on canonical representatives of symmetry classes. But for that, we need to be sure that such canonical representatives always exist.

Of course, there is no free lunch: in the full generality of tcgs, that is not the case⁴.

Example 11. Consider the tcg A , with events, polarities, and causality and follows:



Its symmetry comprises all order-isomorphisms between configurations. The negative symmetry has all order-isomorphisms included in one of the two maximal bijections

$$\begin{array}{ccc} \begin{array}{c} \ominus_1 \\ \vdots \\ \oplus_1 \end{array} & \begin{array}{c} \ominus_2 \\ \vdots \\ \oplus_2 \end{array} & \cong_A^- \\ \begin{array}{c} \ominus_1 \\ \vdots \\ \oplus_2 \end{array} & \begin{array}{c} \ominus_2 \\ \vdots \\ \oplus_1 \end{array} & \end{array}$$

where again, the bijection matches those events in the corresponding position of the diagram. Likewise, the positive symmetry has all order-isomorphisms included in one of:

$$\begin{array}{ccc} \begin{array}{c} \ominus_1 \\ \vdots \\ \oplus_1 \end{array} & \begin{array}{c} \ominus_2 \\ \vdots \\ \oplus_2 \end{array} & \cong_A^+ \\ \begin{array}{c} \ominus_1 \\ \vdots \\ \oplus_2 \end{array} & \begin{array}{c} \ominus_2 \\ \vdots \\ \oplus_1 \end{array} & \end{array}$$

forming, altogether, a tcg. Then, the endosymmetry

$$\begin{array}{ccc} \begin{array}{c} \ominus_1 \\ \vdots \\ \oplus_1 \end{array} & & \begin{array}{c} \ominus_2 \\ \vdots \\ \oplus_1 \end{array} \\ \cong_A & & \end{array}$$

⁴The following example is due to Marc de Visme.

which is neither positive nor negative, uniquely factors as

$$\begin{array}{c} \ominus_1 \quad \ominus_2 \\ \vdots \quad \vdots \\ \oplus_1 \end{array} \cong_A^- \begin{array}{c} \ominus_2 \quad \ominus_1 \\ \vdots \quad \vdots \\ \oplus_2 \end{array} \cong_A^+ \begin{array}{c} \ominus_2 \quad \ominus_1 \\ \vdots \quad \vdots \\ \oplus_1 \end{array}$$

which is not formed of endosymmetries. So this configuration is not canonical, but its only symmetric $\{\ominus_1, \ominus_2, \oplus_2\}$ is not canonical either, for the same reason.

Fortunately, no such pathological example arises in the games that (to our knowledge) have found a use in semantics of logics and programming languages. Next we shall propose the existence of a canonical representative as a new axiom for tcgs, and show that it is preserved by all useful constructions on games.

4.2 Representable games

The axiom of *representability* simply requires the existence of canonical representatives.

Definition 12. *Consider A a tcg.*

*We say that A is **representable** iff for all $\mathbf{x} \in \mathcal{C}_{\cong}(A)$, there is $\underline{\mathbf{x}} \in \mathbf{x}$ canonical.*

If A is representable we may consider fixed in advance a choice, for every symmetry class $\mathbf{x} \in \mathcal{C}_{\cong}(A)$, of a canonical representative $\underline{\mathbf{x}} \in \mathcal{C}(A)$. For this to be a reasonable condition on tcgs, we must check that all the common game constructions preserve representability.

Basic constructions. First, we review the common game constructions that have few interactions with the symmetry. Clearly, the empty game is representable. We have:

Lemma 13. *Consider A, B representable tcgs. Then,*

- (1) A^\perp is representable,
- (2) $A \parallel B$ is representable.

Proof. (1) the dual exchanges \cong_A^+ and \cong_A^- and the definition of canonical is symmetric.

(2) If $\mathbf{x}_A \parallel \mathbf{x}_B \in \mathcal{C}_{\cong}(A \parallel B)$, we simply set $\underline{\mathbf{x}_A \parallel \mathbf{x}_B} = \underline{\mathbf{x}_A} \parallel \underline{\mathbf{x}_B}$. Canonicity follows directly from that of $\underline{\mathbf{x}_A}$ and $\underline{\mathbf{x}_B}$, exploiting the fact that any endosymmetry

$$\theta : \underline{\mathbf{x}_A} \parallel \underline{\mathbf{x}_B} \cong_{A \parallel B} \underline{\mathbf{x}_A} \parallel \underline{\mathbf{x}_B}$$

must have the form $\theta = \theta_A \parallel \theta_B$ for endosymmetries $\theta_A : \underline{\mathbf{x}_A} \cong_A \underline{\mathbf{x}_A}$ and $\theta_B : \underline{\mathbf{x}_B} \cong_B \underline{\mathbf{x}_B}$. \square

The above are the game constructions used in the compact closed structure of thin concurrent games. With similarly direct proofs, we cover all the frequent constructions on tcgs that are essentially independent of symmetry: the shifts $\uparrow A$ (resp. $\downarrow A$) which prefix the game A with a new negative (resp. positive) move (see *e.g.* [7]), the *sum* $\sum_{i \in I} A_i$ having all A_i in pairwise conflict (see *e.g.* [7]), the linear arrow $M \multimap N$ of negative M, N – for all those, preservation of representability is direct. What requires more care is the fact that the constructions that introduce symmetry do indeed preserve representability.

HO exponential. We start with the Hyland-Ong style exponential. Recall that it takes an *arena* in the usual Hyland-Ong sense, *i.e.* a forestial partial order, without symmetry. We refer to [5] for the definition of $!_{HO}A$ for A an arena and the associated notations.

We have the proposition:

Proposition 14. *For A any arena, $!_{HO}A$ is a representable thin concurrent game.*

Proof. Within this proof (and only), by $!A$ we mean $!_{HO}A$. Those configurations $x \in \mathcal{C}(!_{HO}A)$ with exactly one initial move are entirely determined by:

- (1) their label $\text{lbl}(\min(x))$,
- (2) their copy index $\text{ind}(\min(x))$,
- (3) for each $\min(x) \rightarrow a$, the sub-configuration starting with a .

Any $x \in \mathcal{C}(!A)$ with index i , label a and sub-configurations x_1, \dots, x_n may be written

$$x = i \cdot (\{x_1, \dots, x_n\} \multimap a) \in \mathcal{C}(!A)$$

where each x_i is written similarly, with a notation inspired from intersection types. But then, using that similarly any x_j is written $i_j \cdot (X_j \multimap a_j)$, we may rewrite x as

$$i \cdot ((i_1 \cdot x'_1, \dots, i_n \cdot x'_n) \multimap a)$$

where each $x'_j = X_j \multimap a_j$. Going one step further, write

$$x = i \cdot ((i_1^1 \cdot x_1^1, \dots, i_{p_1}^1 \cdot x_{p_1}^1) \multimap \dots \multimap (i_1^m \cdot x_1^m, \dots, i_{p_m}^m \cdot x_{p_m}^m) \multimap a),$$

regrouping sub-trees by symmetry classes. If x is to be canonical, then for any $1 \leq k \leq m$, any i_l^k and $i_{l'}^k$ should be swapped by an endosymmetry; implying $x_l^k = x_{l'}^k$. So we set

$$x' = i \cdot ((i_1^1 \cdot x_1^1, \dots, i_{p_1}^1 \cdot x_1^1) \multimap \dots \multimap (i_1^m \cdot x_1^m, \dots, i_{p_m}^m \cdot x_1^m) \multimap a),$$

which is symmetric to x by construction; moreover if for all $1 \leq k \leq m$, x_1^k is assumed canonical by induction hypothesis, then one may verify that x' is canonical. \square

We omit the details on that last verification, as it is the exact same reasoning as for the AJM exponential, which we give more formally below. Unsurprisingly, the proof for AJM bears much in common with the one above. We started with HO as we believe that the more concrete nature of games obtained through the HO exponential makes the reasoning slightly more transparent: we wish to construct a configuration where any two moves with swappable copy indices have the *exact same* sub-trees below, so that the two copy indices may be simply swapped leaving the remainder of the configuration unchanged.

AJM exponential. The AJM exponential is our main source of non-trivial symmetries.

Lemma 15. *Consider N a representable negative thin concurrent game, i.e. all its minimal events are negative. Then, the thin concurrent game $!N$ is representable.*

Proof. Let $x \in \mathcal{C}(!N)$, of the form $x = \parallel_{i \in I} x_i$, where $x_i \in \mathcal{C}(N)$. Let us partition I as

$$I = \bigsqcup_{k \in K} I_k$$

such that for all $i, j \in I$, $x_i \cong_N x_j$ iff there is some $k \in \mathbb{N}$ such that $i, j \in K$. For each $i \in I$, write $f(i) \in K$ for the corresponding component. For each $k \in K$, fix some $g(k) \in I_k$.

Now, fix $k \in K$. Since N is representable, there is $x_{g(k)} \cong_N \mathbf{canon}(x_{g(k)})$ with $\mathbf{canon}(x_{g(k)})$ canonical. Then for each $j \in I_k$ we replace x_j with $\mathbf{canon}(x_{g(k)})$; or more formally we set

$$x' = \parallel_{i \in I} \mathbf{canon}(x_{g(f(i))}) \in \mathcal{C}(!N).$$

We clearly have $x \cong_{!N} x'$; indeed, for each $i \in I$, we have $x_i \cong_N x_{g(f(i))} \cong_N \mathbf{canon}(x_{g(f(i))})$. Furthermore, x' is canonical. Indeed, writing $x'_i = \mathbf{canon}(x_{g(f(i))})$, consider now any symmetry

$$\theta : \parallel_{i \in I} x'_i \cong_{!N} \parallel_{i \in I} x'_i.$$

By definition, there is $\pi : I \rightarrow I$ a permutation, and for all $i \in I$ a symmetry $\theta_i : x'_i \cong_N x'_{\pi(i)}$. But by construction, this means that we had $x_i \cong_N x_{\pi(i)}$ as well, so $i, \pi(i)$ belong to the same component of the partition and $g(f(i)) = g(f(\pi(i)))$. Therefore, by construction, $x'_i = x'_{\pi(i)}$. But x'_i is canonical, so θ_i decomposes as

$$x'_i \stackrel{\theta_i^-}{\cong}_N^- x'_i \stackrel{\theta_i^+}{\cong}_N^+ x'_i.$$

Setting $\theta^-(i, e) = (\pi(i), \theta_i^-(e))$ and $\theta^+(i, e) = (i, \theta_i^+(e))$, we have the required decomposition of θ , showing that x' is canonical, as required. \square

In the rest of this paper, we aim to make it explicit whenever this condition is required.

5 Quantitative collapse

By now we have added a new condition on games which eliminates some pathological examples, and we have proved that this condition is preserved by all sensible constructions on games. It remains to be seen whether this condition does solve the problem at hand.

5.1 Actions of negative symmetries on strategies

Before we start, recall that for A, B tcgs (which from now on will always be assumed to be representable), $\sigma : A \xrightarrow{S} B$ a strategy and $\mathbf{x}_A \in \mathcal{C}_\cong(A)$, $\mathbf{x}_B \in \mathcal{C}_\cong(B)$, we have set

$$\text{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) = \{x^S \in \mathcal{C}^+(S) \mid x_A^S \cong_A^- \underline{x}_A \ \& \ x_B^S \cong_B^+ \underline{x}_B\},$$

where \underline{x}_A and \underline{x}_B are the canonical representatives given by representability of A and B .

Our next step will be to investigate how negative symmetries act on witnesses. Our starting point for that is the following lemma, Lemma B.4 in [5].

Lemma 16. *Consider $\sigma : S \rightarrow A$ a pre- \sim -strategy, $x^S \in \mathcal{C}(S)$ and $\theta_- : x_A^S \cong_A^- y_A$. Then, there is a unique $\varphi : x^S \cong_S y^S$ s.t. $\sigma\varphi = \theta_+ \circ \theta_- : x_A^S \cong_A y_A^S$ for some $\theta_+ : y_A \cong_A^+ y_A^S$.*

This is our main tool to have negative symmetries act on strategies. If $x^S \in \mathcal{C}(S)$ and $\theta_- : x_A^S \cong_A^- y_A$ presents a change in Opponent's copy indices, we can make θ_- “act on” x^S : Player adapts to the change of Opponent copy indices and presents some $\varphi : x^S \cong_S y^S$.

It is tempting to invoke some group theory here. For any $x \in \mathcal{C}(A)$, we have three groups: the group $\mathcal{S}(x)$ of endosymmetries $\theta : x \cong_A x$, the group $\mathcal{S}_+(x)$ of *positive* endosymmetries, and the group $\mathcal{S}_-(x)$ of *negative* endosymmetries. When applied to symmetry classes, as in $\mathcal{S}(\mathbf{x})$ for $\mathbf{x} \in \mathcal{C}_\cong(A)$, these operations mean $\mathcal{S}(\underline{\mathbf{x}})$. Of course, if $x \cong_A y$ then any $\theta : x \cong_A y$ provides an iso between $\mathcal{S}(x)$ and $\mathcal{S}(y)$ by conjugation. Warning: if $x \cong_A y$ we *do not* necessarily have $\mathcal{S}_-(x)$ and $\mathcal{S}_-(y)$ isomorphic (and of course, likewise for $\mathcal{S}_+(-)$), so the notation $\mathcal{S}_-(\mathbf{x})$ is borderline – we insist that it means $\mathcal{S}_-(\underline{\mathbf{x}})$ and depends on the chosen representative. This shall hopefully cause no confusion.

Now, for $\sigma : S \rightarrow A$ and $x_A \in \mathcal{C}(A)$, it is tempting to make $\mathcal{S}_-(x_A)$ act on the set

$$X = \{x^S \in \mathcal{C}(S) \mid \sigma x^S = x_A\},$$

but for $\theta_- \in \mathcal{S}_-(x_A)$ and $x^S \in X$, there is no reason why the $\varphi : x^S \cong_S y^S$ obtained via Lemma 16 would satisfy $\sigma y^S = x_A$ and hence remain in X .

So we add a bit of wiggling room. For $\mathbf{x}_A \in \mathcal{C}_\cong(A)$, we define the set

$$\sim^\pm \text{wit}^+(\mathbf{x}_A) = \{(x^S, \theta_+) \mid x^S \in \mathcal{C}^+(S), \theta_+ : x_A^S \cong_A^+ \underline{x}_A\}$$

of witnesses for \mathbf{x}_A along with a specific choice of positive symmetry. Then we indeed have:

Proposition 17. *Consider A a tcg and $\mathbf{x}_A \in \mathcal{C}_\cong(A)$. There is a group action*

$$(_ \curvearrowright _) : \mathcal{S}_-(\mathbf{x}_A) \times \sim^\pm \text{wit}^+(\mathbf{x}_A) \rightarrow \sim^\pm \text{wit}^+(\mathbf{x}_A),$$

such that for all $(y^S, \psi_+) = \varphi_- \curvearrowright (x^S, \theta_+)$, there is $\phi^S : x^S \cong_S y^S$ making the diagram

$$\begin{array}{ccc} x_A^S & \xrightarrow{\theta_+} & \underline{x}_A \\ \phi_A^S \downarrow & & \downarrow \varphi_- \\ y_A^S & \xrightarrow{\psi_+} & \underline{x}_A \end{array}$$

commute.

Proof. Consider $(x^S, \theta_+) \in \sim^\pm \text{wit}^+(\mathbf{x}_A)$ and $\varphi_- \in \mathcal{S}_-(\mathbf{x}_A)$. We show that there is unique $\phi^S : x^S \cong_S y^S$ and $\psi_+ : y_A^S \cong_A^+ \underline{\mathbf{x}}_A$ making the following diagram commute:

$$\begin{array}{ccc} x_A^S & \xrightarrow{\theta_+} & \underline{\mathbf{x}}_A \\ \phi_A^S \downarrow & & \downarrow \varphi_- \\ y_A^S & \xrightarrow{\psi_+} & \underline{\mathbf{x}}_A \end{array}$$

For existence, by Lemma 3.19 of [5], $\varphi_- \circ \theta_+ : x_A^S \cong_A \underline{\mathbf{x}}_A$ factors uniquely as

$$\Xi_+ \circ \Xi_- : x_A^S \cong_A \underline{\mathbf{x}}_A.$$

Next, by Lemma 16, there is $\phi^S : x^S \cong_S y^S$ such that we have

$$\phi_A^S = \Omega_+ \circ \Xi_- : x_A^S \cong_A y_A^S$$

for some $\Omega_+ : y_A \cong_A^+ y_A^S$. We then form $\psi_+ = \Xi_+ \circ \Omega_+^{-1}$ to conclude.

For uniqueness, if we have $\varphi_1 : x^S \cong_S y^S$ and $\varphi_2 : x^S \cong_S z^S$ satisfying the requirements,

$$\begin{array}{ccccc} y_A^S & \xrightarrow{(\sigma\varphi_1)^{-1}} & x_A^S & \xrightarrow{\sigma\varphi_2} & z_A^S \\ + \downarrow & & + \downarrow & & \downarrow + \\ \underline{\mathbf{x}}_A & \xrightarrow{\varphi_1^{-1}} & \underline{\mathbf{x}}_A & \xrightarrow{\varphi_-} & \underline{\mathbf{x}}_A \end{array}$$

commutes, so $(\sigma\varphi_2) \circ (\sigma\varphi_1)^{-1} = \sigma(\varphi_2 \circ \varphi_1^{-1})$ is positive, so by Lemma 3.28 of [5] we have $\varphi_2 \circ \varphi_1^{-1} = \text{id}$, so $\varphi_1 = \varphi_2$. \square

Note that in the proof, we have actually not used the representability assumption. However, it will come in to deduce a property useful for elaborate forms of the collapse (namely, in the quantum case). For that, we need the following intermediate lemma.

Lemma 18. *Consider A a representable tcg, $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$, and $x \in \mathcal{C}(A)$ s.t. $x \cong_A^+ \underline{\mathbf{x}}_A$.*

Then, any $\theta : x \cong_A \underline{\mathbf{x}}_A$ factors uniquely as $\theta_- \circ \theta_+$, where $\theta_+ : x \cong_A^+ \underline{\mathbf{x}}_A$ and $\theta_- \in \mathcal{S}_-(\underline{\mathbf{x}}_A)$.

Proof. Fix some $\varphi : x \cong_A^+ \underline{\mathbf{x}}_A$. Now, take $\theta : x \cong_A \underline{\mathbf{x}}_A$. By Lemma 3.19 of [5], θ factors uniquely as $\theta_- \circ \theta_+$, where $\theta_+ : x \cong_A^+ z$ and $\theta_- : z \cong_A^- \underline{\mathbf{x}}_A$ for some $z \in \mathcal{C}(A)$. But then,

$$\varphi \circ \theta^{-1} : \underline{\mathbf{x}}_A \cong_A \underline{\mathbf{x}}_A$$

factors via $(\varphi \circ \theta_+^{-1}) : z \cong_A^+ \underline{\mathbf{x}}_A$ and $\theta_-^{-1} : \underline{\mathbf{x}}_A \cong_A^- z$, so $\underline{\mathbf{x}}_A = z$ follows since $\underline{\mathbf{x}}_A$ is canonical. \square

For A a tcg and $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$, we have previously defined

$$\sim^\pm \text{wit}^+(\mathbf{x}_A) = \{(x^S, \theta_+) \mid x^S \in \mathcal{C}^+(S), \theta_+ : x_A^S \cong_A^+ \underline{\mathbf{x}}_A\}$$

the set of witnesses for $\underline{\mathbf{x}}_A$ up to positive symmetry, *along with* a specific choice of positive symmetry $\theta_+ : x_A^S \cong_A^+ \underline{\mathbf{x}}_A$. We shall now also consider the variation

$$\sim\text{-wit}^+(\mathbf{x}_A) = \{(x^S, \theta) \mid x_A^S \cong_A^+ \underline{\mathbf{x}}_A \text{ \& \; } \theta : x_A^S \cong_A \underline{\mathbf{x}}_A\}$$

where we know that $x_A^S \cong_A^+ \underline{\mathbf{x}}_A$, but $\theta : x_A^S \cong_A \underline{\mathbf{x}}_A$ may not be positive.

Corollary 19. Consider A a representable tcg and $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$. Then, the function

$$\begin{aligned} F &: \sim\text{-}\mathbf{wit}^+(\mathbf{x}_A) \rightarrow \sim^\pm\mathbf{wit}^+(\mathbf{x}_A) \\ (x^S, \theta_- \circ \theta_+) &\mapsto \theta_- \curvearrowright (x^S, \theta_+) \end{aligned}$$

is such that any $X \in \sim^\pm\mathbf{wit}^+(\mathbf{x}_A)$ has exactly $|\mathcal{S}_-(\underline{\mathbf{x}}_A)|$ antecedents.

Proof. The definition of F makes use of the decomposition of all symmetries $\theta : x_A^S \cong_A \underline{\mathbf{x}}_A$ offered by Lemma 18, using canonicity of $\underline{\mathbf{x}}_A$. The statement on the number of antecedents is an immediate consequence of the group action of Proposition 17. \square

The reader might not immediately see the point; in fact we will not use this to establish (2), but it fits in this paper as it is required for more elaborate versions of this construction, in particular in the presence of quantum valuations [7].

5.2 Quantitative synchronization up to symmetry

Let us fix for this section two strategies $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$.

Witnessing strategies and interactions. We write elements of $\sim^\pm\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)$ as triples $(\theta_-^A, x^S, \theta_+^B)$; as an alias for $(x^S, \theta_-^A \parallel \theta_+^B) \in \sim^\pm\mathbf{wit}_\sigma^+(\mathbf{x}_A \parallel \mathbf{x}_B)$. Two witnesses

$$(\theta_-^A, x^S, \theta_+^B) \in \sim^\pm\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B), \quad (\Omega_-^B, x^T, \Omega_+^C) \in \sim^\pm\mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C),$$

are causally compatible iff the composite bijection (see Definition 4) is secured. We write

$$\sim^\pm\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) \bullet \sim^\pm\mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C)$$

for the set of causally compatible pairs $(\mathbf{w}_\sigma, \mathbf{w}_\tau) \in \sim^\pm\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) \times \sim^\pm\mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C)$.

To accompany our notions of witnesses for strategies we shall need to provide witnesses for interactions. If $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$, $\mathbf{x}_B \in \mathcal{C}_{\cong}(B)$ and $\mathbf{x}_C \in \mathcal{C}_{\cong}(C)$, we write

$$\mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C) = \{x^T \otimes x^S \in \mathcal{C}^+(T \otimes S) \mid x_A^S \cong_A^- \underline{\mathbf{x}}_A, \ x_B^S = x_B^T \cong_B \underline{\mathbf{x}}_B, \ \& \ x_C^T \cong_C^+ \underline{\mathbf{x}}_C\}.$$

Like for strategies, we also write $\sim^\pm\mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$ for the set

$$\{(\theta_-^A, x^T \otimes x^S, \theta_+^C) \mid \theta_-^A : x_A^S \cong_A^- \underline{\mathbf{x}}_A, \ x^T \otimes x^S \in \mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C), \ \& \ \theta_+^C : x_C^T \cong_C^+ \underline{\mathbf{x}}_C\}.$$

interaction witnesses along with specific symmetries to the game. Finally, we write:

$$\mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_C) = \{x^T \otimes x^S \in \mathcal{C}^+(T \otimes S) \mid x_A^S \cong_A^- \underline{\mathbf{x}}_A, \ \& \ x_C^T \cong_C^+ \underline{\mathbf{x}}_C\}$$

for the variant of $\mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$ with no constraint in B . Clearly, we have:

Lemma 20. Consider $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$, $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$ and $\mathbf{x}_C \in \mathcal{C}_{\cong}(C)$. Then:

$$\mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_C) = \bigsqcup_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} \mathbf{int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$$

where the notation \sqcup means the plain set-theoretic union when it is disjoint.

Proof. Simply partition interactions according to the symmetry class reached in B . \square

Interactions up to symmetry. We start with a more explicit variant of Lemma 6.

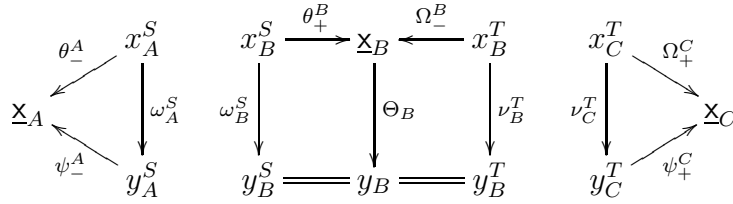
Lemma 21. *For any pair of causally compatible witnesses*

$$(\theta_-^A, x^S, \theta_+^B) \in \sim^\perp \text{-wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B), \quad (\Omega_-^B, x^T, \Omega_+^C) \in \sim^\perp \text{-wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C),$$

there are unique symmetries $\omega^S : x^S \cong_S y^S, \nu^T : x^T \cong_T y^T$, $\Theta_B : \underline{\mathbf{x}}_B \cong_B y_B$ and witness

$$(\psi_-^A, y^T \otimes y^S, \psi_+^C) \in \sim^\perp \text{-int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$$

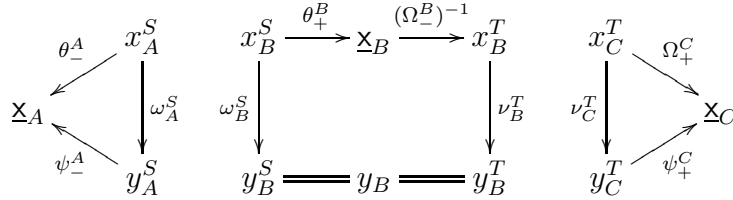
with $y_B^S = y_B^T = y_B$, such that the following diagrams commute:



Proof. By Lemma 6, there are unique symmetries $\omega^S : x^S \cong_S y^S, \nu^T : x^T \cong_T y^T$, and

$$(\psi_-^A, y^T \otimes y^S, \psi_+^C) \in \sim^\perp \text{-int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$$

with $y_B^S = y_B^T = y_B$, such that the following diagrams commute:



We simply set $\Theta_B : \underline{\mathbf{x}}_B \rightarrow y_B$ as either path around the center diagram. □

Thanks to the previous section we may reverse this operation, as shown below.

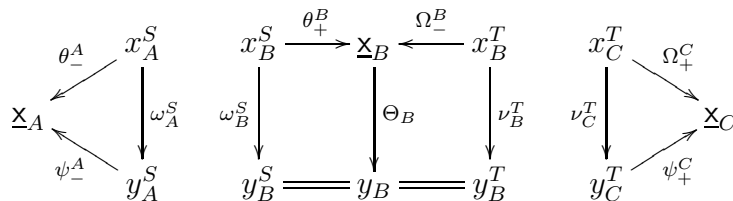
Lemma 22. *For any symmetry $\Theta_B : \underline{\mathbf{x}}_B \cong_B y_B$ and any witness*

$$(\psi_-^A, y^T \otimes y^S, \psi_+^C) \in \sim^\perp \text{-int}_{\tau \otimes \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)$$

with $y_B^S = y_B^T = y_B$, there are unique symmetries $\omega^S : x^S \cong_S y^S, \nu^T : x^T \cong_T y^T$ and

$$(\theta_-^A, x^S, \theta_+^B) \in \sim^\perp \text{-wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B), \quad (\Omega_-^B, x^T, \Omega_+^C) \in \sim^\perp \text{-wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C),$$

a pair of causally compatible witnesses, such that the following diagrams commute:



Proof. The first step is to factor Θ_B^{-1} in two ways, as in the diagram

$$\begin{array}{ccccc}
 & & z_B^1 & \xrightarrow{\Phi_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 & & \\
 & & \swarrow & & \uparrow & & \searrow & & \\
 \underline{x}_A & \xleftarrow{\psi_-^A} & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T & \xrightarrow{\psi_+^C} & \underline{x}_C
 \end{array}$$

following Lemma 3.19 of [5]. By Lemma 16 we can make Φ_-^B act on σ . This yields

$$\lambda_-^A : x_A^S \cong_A^- y_A^S, \quad \omega^S : x^S \cong_S y^S, \quad \Delta_+^B : x_B^S \cong_B^+ z_B^1,$$

unique such that the following diagram commutes:

$$\begin{array}{ccccc}
 & & x_B^S & \xrightarrow{\Delta_+^B} & z_B^1 & \xrightarrow{\Phi_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 & & \\
 & & \downarrow \omega_B^S & & \swarrow & & \uparrow & & \searrow & & \\
 \underline{x}_A & \xleftarrow{\psi_-^A} & y_A^S & \xlongequal{\quad} & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T & \xrightarrow{\psi_+^C} & \underline{x}_C
 \end{array}$$

leaving in grey the irrelevant parts of the full diagram for context. Setting $\theta_-^A = \psi_-^A \circ \lambda_-^A$ and $\theta_+^B = \Phi_+^B \circ \Delta_+^B$, we have found data making the following diagram commute:

$$\begin{array}{ccccc}
 & & x_A^S & \xrightarrow{\theta_-^A} & \underline{x}_A & & x_B^S & \xrightarrow{\theta_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 & & \\
 & & \downarrow \omega_A^S & & \swarrow & & \downarrow \omega_B^S & & \swarrow & & \uparrow & & \searrow & & \\
 \underline{x}_A & \xleftarrow{\psi_-^A} & y_A^S & \xlongequal{\quad} & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T & \xrightarrow{\psi_+^C} & \underline{x}_C
 \end{array}$$

We shall now prove uniqueness of this data. Assume that we have other symmetries

$$\gamma_-^A : u_A^S \cong_A^- \underline{x}_A, \quad \varpi^S : u^S \cong_S y^S, \quad \gamma_+^B : u_B^S \cong_B^+ \underline{x}_B,$$

making the following diagram commute:

$$\begin{array}{ccccc}
 & & u_A^S & \xrightarrow{\gamma_-^A} & \underline{x}_A & & u_B^S & \xrightarrow{\gamma_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 & & \\
 & & \downarrow \varpi_A^S & & \swarrow & & \downarrow \varpi_B^S & & \swarrow & & \uparrow & & \searrow & & \\
 \underline{x}_A & \xleftarrow{\psi_-^A} & y_A^S & \xlongequal{\quad} & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T & \xrightarrow{\psi_+^C} & \underline{x}_C
 \end{array}$$

Then, it follows that the following diagram also commutes:

$$\begin{array}{ccccc}
 & & u_B^S & \xrightarrow{(\Phi_+^B)^{-1} \circ \gamma_+^B} & z_B^1 & \xrightarrow{\Phi_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 & & \\
 & & \downarrow \varpi_B^S & & \swarrow & & \uparrow & & \searrow & & \\
 \underline{x}_A & \xleftarrow{\psi_-^A} & y_A^S & \xlongequal{\quad} & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T & \xrightarrow{\psi_+^C} & \underline{x}_C
 \end{array}$$

By uniqueness for Lemma 16, it follows that $u^S = x^S$, $\omega^S = \varpi^S$, $\lambda_-^A = (\psi_-^A)^{-1} \circ \gamma_-^A$ so $\gamma_-^A = \theta_-^A$, and $(\Phi_+^B)^{-1} \circ \gamma_+^B = \Delta_+^B$ so $\gamma_+^B = \theta_+^B$. Altogether, we have proved that there are

$$\theta_-^A : x_A^S \cong_S \underline{x}_A, \quad \omega^S : x^S \cong_S y^S, \quad \theta_+^B : x_B^S \cong_B^+ \underline{x}_B,$$

unique making the following diagram commutes:

$$\begin{array}{ccccccc} & & x_A^S & & x_B^S & \xrightarrow{\theta_+^B} & \underline{x}_B & \xleftarrow{\Psi_-^B} & z_B^2 \\ & \swarrow \theta_-^A & \downarrow \omega_A^S & & \downarrow \omega_B^S & & \downarrow \Theta_B & \nearrow \Psi_+^B & \\ \underline{x}_A & & y_A^S & & y_B^S & \xlongequal{\quad} & y^B & \xlongequal{\quad} & y_C^T & \nearrow \psi_+^C & \underline{x}_C \\ & \nwarrow \psi_-^A & & & & & & & & & \end{array}$$

The lemma follows by performing the exact same reasoning on the right hand side. \square

5.3 Witnesses of interaction

With Lemmas 21 and 22 we have done the hardest part of the job; but to collect the fruits of that work we need to introduce some additional notation. We write $\mathcal{S}(\underline{x}_B)$ for the set of endosymmetries on \underline{x}_B . Let us fix a choice, for every $x \in \underline{x}_B$, of some

$$\kappa_x : x \cong_B \underline{x}_B.$$

Transporting through $(\kappa_x)_{x \in \underline{x}_B}$ gives a bijection, for any two $x, y \in \underline{x}_B$, between the set of symmetries $x \cong_B y$ and the set $\mathcal{S}(\underline{x}_B)$. If $\theta \in \mathcal{S}(\underline{x}_B)$ and $x, y \in \underline{x}_B$, let us write

$$\theta[x, y] : x \cong_B y$$

the transported symmetry obtained as $\kappa_y^{-1} \circ \theta \circ \kappa_x$.

Corollary 23. *There is a bijection*

$$\Upsilon : \quad \sim^\perp \text{wit}_\sigma^+(\underline{x}_A, \underline{x}_B) \bullet \sim^\perp \text{wit}_\tau^+(\underline{x}_B, \underline{x}_C) \quad \rightarrow \quad \sim^\perp \text{int}_{\tau \oplus \sigma}^+(\underline{x}_A, \underline{x}_B, \underline{x}_C) \times \mathcal{S}(\underline{x}_B)$$

such that for every pair of causally compatible witnesses

$$(\mathbf{w}_1, \mathbf{w}_2) = ((\theta_-^A, x^S, \theta_+^B), (\Omega_-^B, x^T, \Omega_+^C)) \in \sim^\perp \text{wit}_\sigma^+(\underline{x}_A, \underline{x}_B) \bullet \sim^\perp \text{wit}_\tau^+(\underline{x}_B, \underline{x}_C),$$

writing $((\psi_-^A, y^T \otimes y^S, \psi_+^C), \varphi) = \Upsilon(\mathbf{w}_1, \mathbf{w}_2)$, $y_B = y_B^S = y_B^T$, $\Theta_B = \varphi[\underline{x}_B, y_B]$, there are $\omega^S : x^S \cong_S y^T$ and $\nu^T : x^T \cong_T y^T$ such that the following diagrams commute:

$$\begin{array}{ccccccc} & & x_A^S & & x_B^S & \xrightarrow{\theta_+^B} & \underline{x}_B & \xleftarrow{\Omega_-^B} & x_B^T & & x_C^T \\ & \swarrow \theta_-^A & \downarrow \omega_A^S & & \downarrow \omega_B^S & & \downarrow \Theta_B & & \downarrow \nu_B^T & & \downarrow \nu_C^T \\ \underline{x}_A & & y_A^S & & y_B^S & \xlongequal{\quad} & y_B & \xlongequal{\quad} & y_B^T & & y_C^T \\ & \nwarrow \psi_-^A & & & & & & & & \nearrow \psi_+^C & \underline{x}_C \end{array}$$

Proof. Straightforward from Lemmas 21 and 22. Note that ω^S and ν^T are unique; the requirements of the diagrams constrain them entirely due to local injectivity of σ, τ . \square

The commutation of this diagram is required for situations where one would exploit this in the presence of valuations on configurations that are *typed* and transported coherently through symmetry, such as for quantum valuations [7]. However, if one is merely interested in *counting* the witnesses, then the take home message is:

Corollary 24. *For any $\mathbf{x}_A \in \mathcal{C}_{\cong}(A)$, $\mathbf{x}_B \in \mathcal{C}_{\cong}(B)$ and $\mathbf{x}_C \in \mathcal{C}_{\cong}(C)$, we have*

$$|\sim^{\pm}\text{wit}_{\sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B) \bullet \sim^{\pm}\text{wit}_{\tau}^{+}(\mathbf{x}_B, \mathbf{x}_C)| = |\sim^{\pm}\text{int}_{\tau \circ \sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)| \times |\mathcal{S}(\mathbf{x}_B)|. \quad (4)$$

If we know that the strategies to be composed do not deadlock, then this can be simplified further.

Corollary 25. *Assume $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ do not deadlock. Then,*

$$|\sim^{\pm}\text{wit}_{\sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B)| \times |\sim^{\pm}\text{wit}_{\tau}^{+}(\mathbf{x}_B, \mathbf{x}_C)| = |\sim^{\pm}\text{int}_{\tau \circ \sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)| \times |\mathcal{S}(\mathbf{x}_B)|,$$

Proof. By hypothesis, causal compatibility is always satisfied. Therefore,

$$\sim^{\pm}\text{wit}_{\sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B) \bullet \sim^{\pm}\text{wit}_{\tau}^{+}(\mathbf{x}_B, \mathbf{x}_C) = \sim^{\pm}\text{wit}_{\sigma}^{+}(\mathbf{x}_A, \mathbf{x}_B) \times \sim^{\pm}\text{wit}_{\tau}^{+}(\mathbf{x}_B, \mathbf{x}_C)$$

and the result follows from Corollary 24. \square

This takes us close to Equation 2. One may wonder what is left to conclude; a hint is the fact that for now, in this section, we have *not used canonicity of representatives*.

5.4 Witnesses and canonicity

The moral of Equation 4 seems clear: on the left hand side witnesses have the liberty to pick any positive symmetry on respectively B and B^{\perp} to interact, whereas on the right hand side they must match on the nose. Adding $|\mathcal{S}(\mathbf{x}_B)|$ on the right balances this out.

Let us look deeper into this. From now on, we will rely heavily on canonicity of representatives. A first consequence of that is the following:

Lemma 26. *If B is representable, then for all $\mathbf{x}_B \in \mathcal{C}_{\cong}(B)$, we have*

$$|\mathcal{S}(\mathbf{x}_B)| = |\mathcal{S}_{-}(\mathbf{x}_B)| \times |\mathcal{S}_{+}(\mathbf{x}_B)|.$$

Proof. Obvious consequence of the definition of canonicity. \square

Indeed this is almost the definition of canonicity, which states that every endosymmetry on $\underline{\mathbf{x}}_B$ factors uniquely as the composition of a positive and a negative endosymmetries of $\underline{\mathbf{x}}_B$. Almost as obvious is the following fact:

Lemma 27. For any $\sigma : A \xrightarrow{S} B, \tau : B \xrightarrow{T} C, \mathbf{x}_A \in \mathcal{C}_{\cong}(A), \mathbf{x}_B \in \mathcal{C}_{\cong}(B)$, and $\mathbf{x}_C \in \mathcal{C}_{\cong}(C)$,

$$\begin{aligned} |\sim^\pm \mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)| &= |\mathcal{S}_-(\mathbf{x}_A)| \times |\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)| \times |\mathcal{S}_+(\mathbf{x}_B)| \\ |\sim^\pm \mathbf{int}_{\tau \oplus \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)| &= |\mathcal{S}_-(\mathbf{x}_A)| \times |\mathbf{int}_{\tau \oplus \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)| \times |\mathcal{S}_+(\mathbf{x}_C)| \end{aligned}$$

Proof. We only detail the first equality, the reasoning for the other is identical. Let us choose, for every $x \in \mathbf{x}_B$ such that $x \cong_B^+ \underline{x}_B$, some positive symmetry $\kappa_x^B : \underline{x}_B \cong_B^+ x$. Likewise we choose, for each $y \in \mathbf{x}_A$ such that $y \cong_A^- \underline{x}_A$, some $\kappa_y^A : \underline{x}_A \cong_A^- y$.

Now, we form the function:

$$\begin{aligned} G : \sim^\pm \mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) &\rightarrow \mathcal{S}_-(\mathbf{x}_A) \times \mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B) \times \mathcal{S}_+(\mathbf{x}_B) \\ (\theta_-^A, x^S, \theta_+^B) &\mapsto (\theta_-^A \circ \kappa_{x^S}^A, x^S, \theta_+^B \circ \kappa_{x^S}^B) \end{aligned}$$

which is clearly a bijection as positive and negative symmetries are invertible. \square

5.5 Wrapping up

Finally, we are now in position to prove:

Theorem 28. Consider $\sigma : A \xrightarrow{S} B$ and $\tau : B \xrightarrow{T} C$ that do not deadlock, and assume that B is representable. Then, for all $\mathbf{x}_A \in \mathcal{C}_{\cong}(A), \mathbf{x}_C \in \mathcal{C}_{\cong}(C)$, we have

$$(\int (\tau \odot \sigma))_{\mathbf{x}_A, \mathbf{x}_C} = \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} (\int \sigma)_{\mathbf{x}_A, \mathbf{x}_B} \cdot (\int \tau)_{\mathbf{x}_B, \mathbf{x}_C}.$$

Proof. We calculate

$$(\int (\tau \odot \sigma))_{\mathbf{x}_A, \mathbf{x}_C} = |\mathbf{wit}_{\tau \odot \sigma}^+(\mathbf{x}_A, \mathbf{x}_C)| \quad (5)$$

$$= |\mathbf{int}_{\tau \oplus \sigma}^+(\mathbf{x}_A, \mathbf{x}_C)| \quad (6)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} |\mathbf{int}_{\tau \oplus \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)| \quad (7)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} \frac{|\sim^\pm \mathbf{int}_{\tau \oplus \sigma}^+(\mathbf{x}_A, \mathbf{x}_B, \mathbf{x}_C)|}{|\mathcal{S}_-(A)| \cdot |\mathcal{S}_+(C)|} \quad (8)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} \frac{|\sim^\pm \mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)| \cdot |\sim^\pm \mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C)|}{|\mathcal{S}_-(A)| \cdot |\mathcal{S}_-(B)| \cdot |\mathcal{S}_+(C)|} \quad (9)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} \frac{|\mathcal{S}_+(\mathbf{x}_B)| \cdot |\mathcal{S}_-(\mathbf{x}_B)|}{|\mathcal{S}_-(\mathbf{x}_B)|} \cdot |\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)| \cdot |\mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C)| \quad (10)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} |\mathbf{wit}_\sigma^+(\mathbf{x}_A, \mathbf{x}_B)| \cdot |\mathbf{wit}_\tau^+(\mathbf{x}_B, \mathbf{x}_C)| \quad (11)$$

$$= \sum_{\mathbf{x}_B \in \mathcal{C}_{\cong}(B)} (\int \sigma)_{\mathbf{x}_A, \mathbf{x}_B} \times (\int \tau)_{\mathbf{x}_B, \mathbf{x}_C} \quad (12)$$

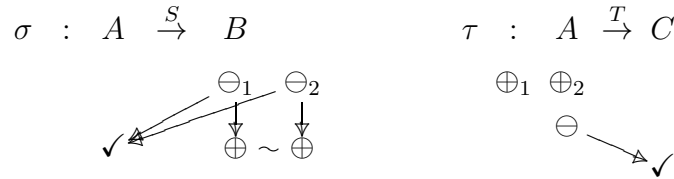
where (5) is by definition, (6) is by Lemma 3, (7) is by Lemma 20, (8) is by Lemma 27, (9) is by Corollary 25, (10) is by Lemma 27 again, (11) is by Lemma 26 exploiting that B is representable; and (12) is by definition. \square

This concludes the proof of (2).

6 Epilogue: back to symmetry classes

To conclude, we show that our original notion of witness based on symmetry classes rather than canonical representatives was, in fact, wrong. We give two counter-examples: the first is geared towards simplicity, while the second aims to bring the counter-example as close as possible to usual models of programming languages. The two examples are, however, powered by the same phenomenon.

Example 29. Consider the games $A^\perp, C = \checkmark$ formed of only one positive move. The game $B = \ominus_1 \ominus_2 \oplus$ has three moves, with \ominus_1 and \ominus_2 symmetric. We consider two strategies:



These are indeed valid strategies in the sense of [5]. Their composition is:

$$\tau \odot \sigma : A \xrightarrow{T \odot S} C$$

$$\checkmark \qquad \checkmark \sim \checkmark$$

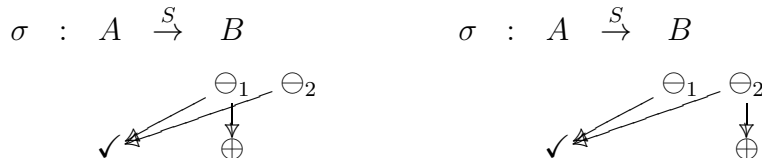
In particular, the non-deterministic choice on the right hand side originates from the choice by σ : to which \ominus_i should it react? In particular,

$$|\text{wit}_{\tau \odot \sigma}(\{\checkmark\}, \{\checkmark\})| = 2,$$

as the two occurrences of \checkmark on the right are not symmetric (this boils down to the fact that \oplus_1 and \oplus_2 cannot be symmetric in τ , by thinness). On the other hand, the only symmetry class of B on which these two may interact is $\{\ominus_1, \ominus_2, \oplus\}$. And we have:

$$|\text{wit}_\sigma(\{\checkmark\}, \{\ominus_1, \ominus_2, \oplus\})| = 1 \qquad |\text{wit}_\tau(\{\oplus_1, \oplus_2, \ominus\}, \{\checkmark\})| = 1$$

In particular, the two configurations of σ responsible for the non-deterministic choice



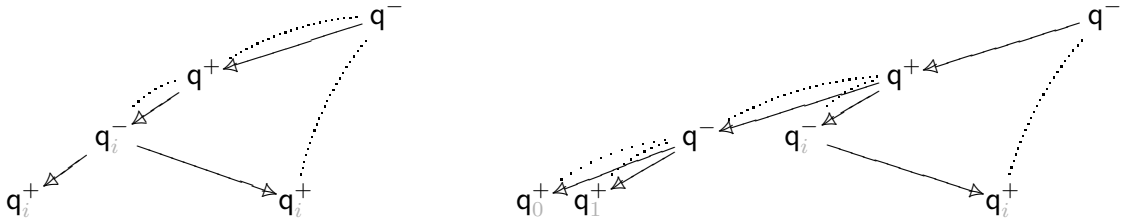
are symmetric, so they form only one symmetry class and are counted only once in $\text{wit}_\sigma(\{\checkmark\}, \{\ominus_1, \ominus_2, \oplus\})$ – whereas they are two distinct elements of $\text{wit}_\sigma^+(\{\checkmark\}, \{\ominus_1, \ominus_2, \oplus\})$.

This also shows that it is *not* the case that configurations in $\text{wit}_\sigma^+(x)$ are canonical representatives of symmetry classes – they are better than that, as they get it right where symmetry classes get it wrong.

We now show essentially the same example in a more “programming language” style.

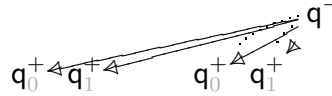
Example 30. Consider a basic game o , with a unique move q^- . Consider strategies:

$$\sigma : !o \xrightarrow{S} (!o \multimap o) \multimap !o \multimap o \quad \tau : ((!o \multimap o) \multimap !o \multimap o) \xrightarrow{T} !o \multimap o$$



Note that the moves on the left are only there to ensure the $+$ -covered hypothesis. Their composition is:

$$\tau \odot \sigma : !o \xrightarrow{T \odot S} !o \multimap o$$



Now, as for the previous example, we observe:

$$|\text{wit}_{\tau \odot \sigma}([q, q], [q] \multimap q)| = 2$$

using an intersection type like notation for symmetry classes on the game, which hopefully is clear. The reader may check that there is a unique symmetry class on $(!o \multimap o) \multimap !o \multimap o$ on which the strategies may match to produce this via $+$ -covered configurations, namely

$$([q, q] \multimap q) \multimap [q] \multimap q$$

in the same intersection-type like notation. And we have

$$|\text{wit}_\sigma([q, q], ([q, q] \multimap q) \multimap [q] \multimap q)| = 1 \quad |\text{wit}_\tau((([q, q] \multimap q) \multimap [q] \multimap q), ([q], q))| = 1$$

for the same reason as in the previous example.

These strategies are not quite terms but they are very well-behaved, in particular visible and parallel innocent. This counter-example does not quite contradict the claims of [3] because there strategies are more constrained (in particular they are well-bracketed) and the positions of interest (matching the points of the web) are complete. It is plausible that this makes this pathology disappear – in particular Example 30 exploits non-well bracketed behaviour, but this is pure speculation. In any case concrete witnesses as developped here are definitely better behaved, and are recommended in all situations.

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